Surface Energy Budget

Please read Bonan Chapter 13

Energy Budget Concept

- For the land surface,
 - Energy in = Radiation
 - Energy Out = Radiation + Turbulent fluxes of "sensible" and "latent" heat
 - Change in energy = changes in temperature of soil, plants, water, and air

Energy Budget Concept

For any system,

(Energy in) – (Energy out) = (Change in energy)

- For the land surface,
 - Energy in =?
 - Energy Out =?
 - Change in energy = ?

Surface Radiation Budget

- Shortwave
 - Down (solar constant, seasonal and diurnal geometry, atmospheric attenuation, clouds and aerosol)
 - Up (albedo)
- Longwave
 - Down (emission from atmosphere depends on temperature profile, water vapor, clouds)
 - Up (surface temperature, emissivity)

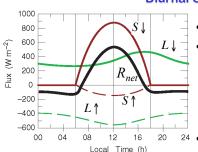
Surface Radiation Budget

$$R_{net} = SW \downarrow -SW \uparrow + LW \downarrow -LW \uparrow$$
$$= SW \downarrow (1 - albedo) + LW \downarrow (1 - \varepsilon) + \varepsilon \sigma T_s^4$$

- Shortwave
 - Down (solar constant * cosZ * transmissivityairmass)
 - Up (albedo * SW down)
- Longwave
 - Down (complicated! Weighted average of σT_a⁴)
 - Up ($\varepsilon \sigma T_s^4$)

Radiation Budget Components

Diurnal Cycle



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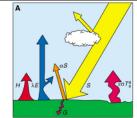
- Net solar follows cos(z)
- LW fluxes much less variable (εσT⁴)
 - LW up follows surface T as it warms through day
 - LW down changes little
 - LW net opposes SW_{net}
- R_{net} positive during day, slightly negative at night

heat

flux

Land Surface Energy Budget

- Very little of the energy gained by net radiation is stored in the ground (G)
- Most is emitted as LW IR and turbulent fluxes of sensible (H) and latent heat (LE)
- Latent energy is then released into atmosphere when vapor condenses



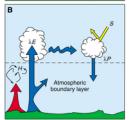


Fig. 1. Interactions between the land surface and the atmosphere that have direct impacts on the physical climate system. (A) Surface radiation budget. (B) Effect of heat fluxes on the atmosphere.

Surface Energy Budget

Storage change = Energy in - energy out

$$\rho c \frac{\Delta T}{\Delta t} \Delta z = (S \downarrow -S \uparrow + L \downarrow -L \uparrow) - \frac{H}{H} + \lambda E = G$$

$$R_{net} = (S \downarrow -S \uparrow) + (L \downarrow -L \uparrow) = H + \lambda E + G$$

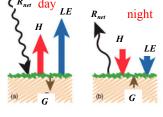
$$\underset{net \ SW}{\text{net \ LW}} \qquad \underset{\text{sensible}}{\text{heat flux}} \qquad \underset{\text{evap}}{\text{latent heat}}$$

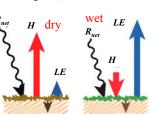
$$\underset{\text{of vaporization}}{\text{of vaporization}} \qquad \underset{\text{ground}}{\text{ground}}$$

Role of the land surface:

Partition of net radiation into turbulent fluxes & storage

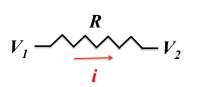
Surface Energy Budgets

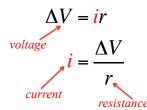




- R_{net} = H + LE + G ~ H + LE
- Daytime turbulent fluxes upward
- Night: turbulent fluxes downward (dew or frost!)
- Dry surfacesR_{net} ~ H
- Wet surfaces
 R_{net} ~ LE

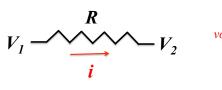
Remember Ohm's Law?

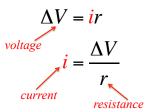




- Flow of current through a resistor is ratio of difference in potential to resistance (this is just a "definition of resistance")
- This is another form of our familiar concept of stuff flowing from high concentration to low concentration (like "Fickian Diffusion")

Heat Fluxes ~ Currents

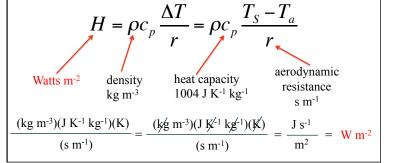




- Sensible heat flux
 - Driving potential is a difference in temperature
 - H is proportional to △T
- Latent heat flux
 - Driving potential is a difference in vapor pressure
 - LE is proportional to ∆e

Sensible Heat Flux

- Driving potential is a difference in temperature
- H is proportional to ΔT

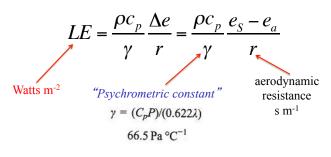


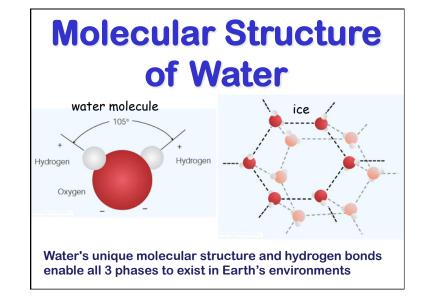
Water Vapor Pressure

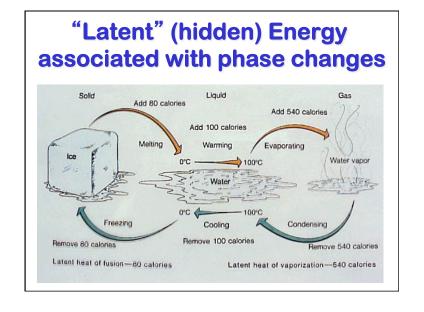
- Molecules in an air parcel all contribute to pressure
- Each subset of molecules (e.g., N₂, O₂, H₂O) exerts a partial pressure
- The VAPOR PRESSURE (e), is the pressure exerted by water vapor molecules in the air

Latent Heat Flux

- Driving potential is a difference in water vapor pressure
- LE is proportional to ∆e

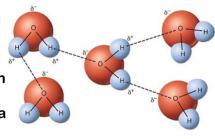






Why does it take so much energy to evaporate water?

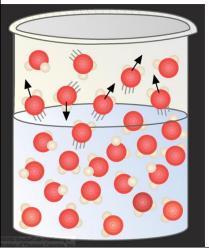
- In the liquid state, adjacent water molecules attract one another
- This same hydrogen bond accounts for surface tension on a free water surface



"plus" charge on hydrogen in one water molecule attracts the "minus" charge on a neighbor's oxygen

Evaporation must break these hydrogen bonds

Water Vapor Saturation

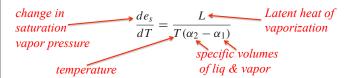


- Water molecules move between the liquid and gas phases
- When the rate of water molecules entering the liquid equals the rate leaving the liquid, we have equilibrium
- The air is said to be saturated with water vapor at this point

Clausius-Clapeyron Eqn

(see Monteith & Unsworth, pp 11-13)

• From Second Law of Thermodynamics



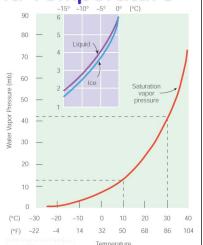
Approximate but very accurate 0° to 35° C

$$e_s(T) = e_s(T^*) \exp\{A(T - T^*)/(T - T')\}$$

where A = 17.27, $T^* = 273$ K ($e_s(T^*) = 0.611$ kPa), and T' = 36 K.

Saturation and Temperature

- The saturation vapor pressure of water increases with temperature
 - At higher T, faster water molecules in liquid escape more frequently causing equilibrium water vapor concentration to rise
 - We sometimes say "warmer air can hold more water"
- There is also a vapor pressure of water over an ice surface
 - The saturation vapor pressure above solid ice is less than above liquid water



Latent Heat Flux

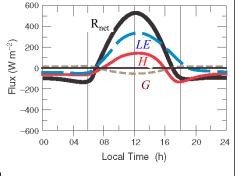
- Driven by difference in vapor pressure
- Over open water e_{surface} = e_{sat}(T_s)
- Over vegetation, liquid water is evaporating inside tiny openings in leaves called *stomata* (singular "stomate")
- Evapotranspiration = latent heat flux is driven by the vapor pressure deficit

$$vpd = (e_{sat}(T_s) - e_a)$$

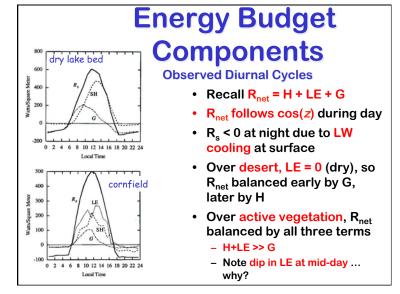
$$LE = \frac{\rho c_p}{\gamma} \frac{e_s - e_a}{r} = \frac{\rho c_p}{\gamma} \frac{e_{sat}(T_s) - e_a}{r}$$

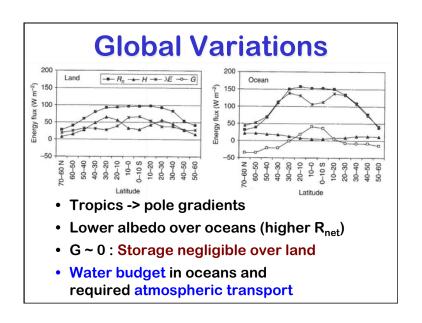
Idealized Diurnal Cycle R_{net} follows

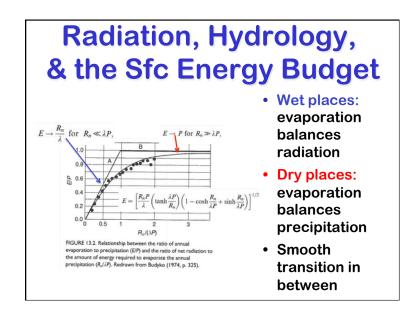
- R_{net} follows cos(z) during day, negative at night (LW cooling)
- Downward turbulent fluxes at night
- Ground heat flux smaller: downward durng day and up at night

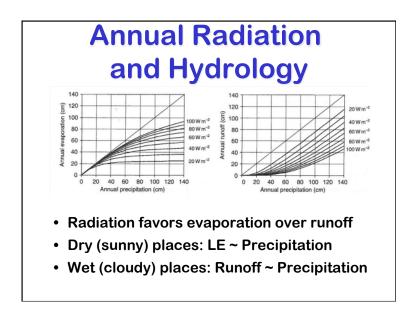


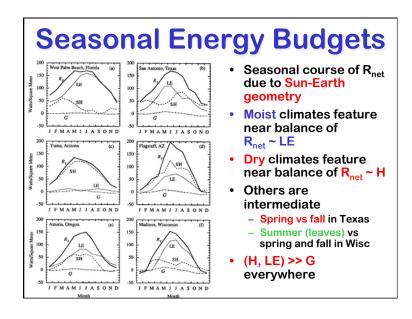
 $R_{net} = H + LE + G$ ~ H + LE

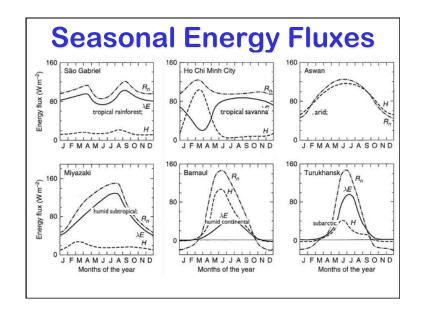


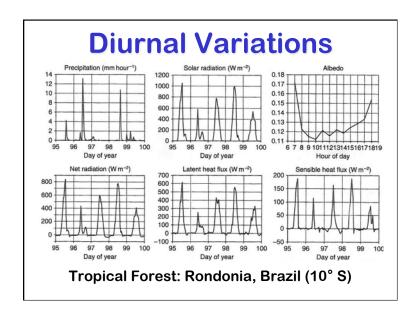


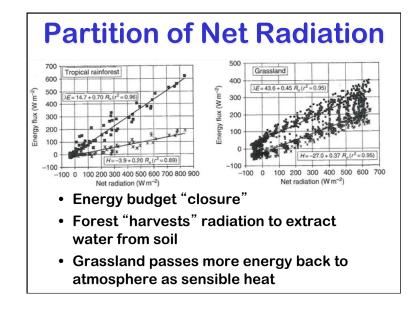


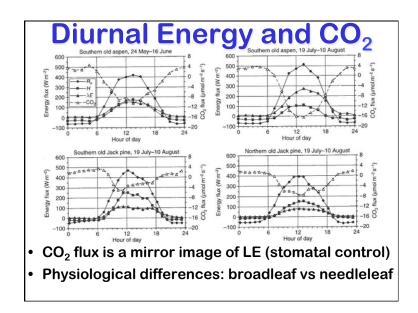


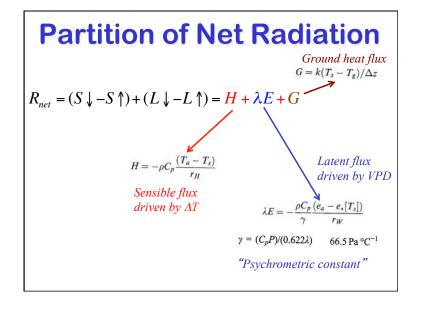












9

Surface Energy Budget

Energy in = energy out + storage change

$$(1-r)S\downarrow +\varepsilon L\downarrow =\varepsilon \sigma (T_s+273.15)^4 + H + \lambda E + G$$

$$(1-r)S\downarrow +\varepsilon L\downarrow =\varepsilon \sigma (T_s+273.15)^4 -\rho C_p \frac{(T_a-T_s)}{r_H} -\frac{\rho C_p}{\gamma} \frac{(e_a-e_*[T_s])}{r_W} + k \frac{(T_s-T_g)}{\Delta z}$$

- Can "solve" for surface temperature
- Physical properties: albedo, emissivity, heat capacity, soil conductivity & temperature
- "Resistances" are properties of the turbulence ... depend sensitively on H!

Clausius-Clapeyron Eqn

(see Monteith & Unsworth, pp 11-13)

Approximate but very accurate 0° to 35° C

$$e_{\rm S}(T) = e_{\rm S}(T^*) \exp\{A(T-T^*)/(T-T')\}$$
 where $A=17.27, T^*=273$ K ($e_{\rm S}(T^*)=0.611$ kPa), and $T'=36$ K.

• Slope $s = \Delta(e_s)/\Delta T$ accurate from 0° to 40° C $= \lambda M_{\rm w} e_{\rm s}(T)/(RT^2)$

 $\lambda = L = 2.48 \text{ J kg}^{-1}$ $M_W \text{ (mol wt water)} = 0.018 \text{ mol kg}^{-1}$

 $R = 8.314 \text{ J mol}^{-1} \text{ K}^{-1}$ ("universal gas constant")

Penman-Monteith Equation

"Thermodynamic" energy balance

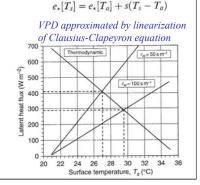
$$\lambda E = (R_n - G) - H = (R_n - G) + \rho C_p (T_a - T_s) / r_H$$

"Turbulent" energy balance

$$\lambda E = \frac{\rho C_p}{\gamma} \frac{(e_*[T_a] + s(T_s - T_a) - e_a}{r_W}$$

Solve for surface temperature

$$T_s - T_a = (r_H/\rho C_p)(R_n - G - \lambda E)$$



Solutions to P-M Equation

Latent heat flux

$$\lambda E = \frac{s(R_n - G) + \rho C_\rho(e_*[T_a] - e_a)/r_H}{s + \gamma(r_W/r_H)}$$

Sensible heat flux

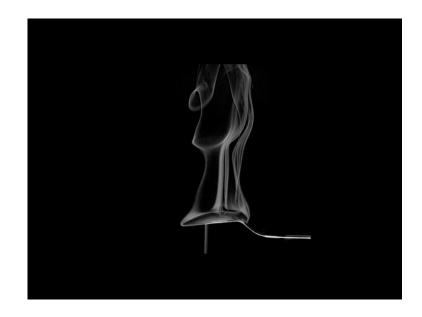
$$H = \frac{(R_n - G)\gamma^* - \rho C_p(e_*[T_a] - e_a)/r_H}{s + \gamma^*}$$

Surface temperature

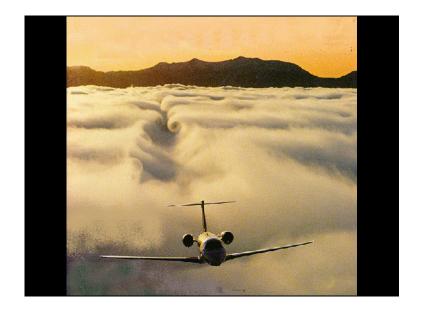
$$T_s = T_a + \frac{(R_n - G)\gamma^* r_H / \rho C_p - (e_*[T_a] - e_a)}{s + \gamma^*}$$

Turbulent Fluxes

Please read Bonan, Chapter 14







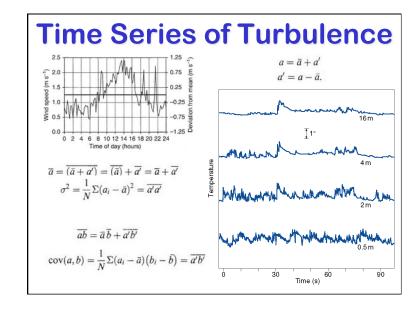
Richardson's Rhyme

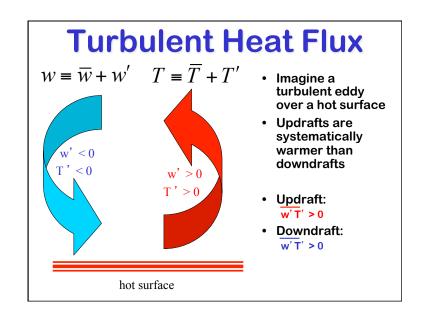
- "Big whorls have little whorls, Which feed on their velocity; And little whorls have lesser whorls, And so on to viscosity"
 - Lewis Richardson, The supply of energy from and to Atmospheric Eddies 1920
- "Great fleas have little fleas Upon their backs to bite 'em, And little fleas have lesser fleas, And so, ad infinitum"
 - Augustus De Morgan (19th century mathematician, parodying Jonathon Swift, 1733)

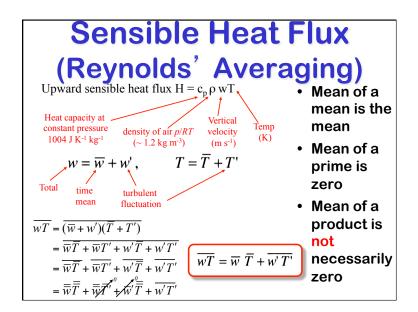


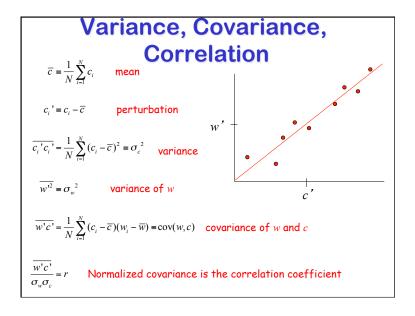
Sonic Anemometer

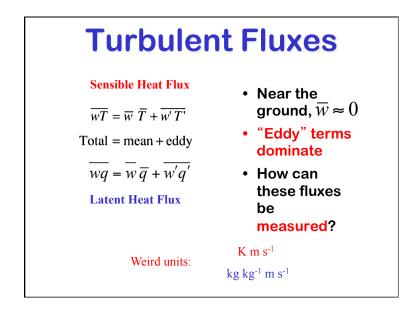
- Measures elapsed time for sound pulses to cross air in 3D
- Speed of sound is a known function of temperature
- Relative motion determined accurately in 3D
- Very fast instrument response time

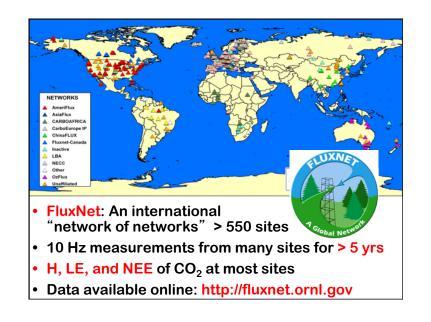












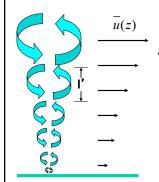








Surface-Layer Mixing



- Turbulent eddies near the surface act to mix atmospheric properties (T, q, u) and reduce vertical gradients
- Assume a characteristic length scale *l'* for eddy mixing, then

$$u = u(z) + u'$$

$$u' = -l' \frac{\partial \overline{u}}{\partial z}$$

If eddies are isotropic

(length and depth similar), then

$$|w'| \sim |u'|$$
, so $w' \sim l' \frac{\partial \overline{u}}{\partial z}$

Surface Layer Stress

$$\tau_{x} = -\rho \overline{w'u'}$$

$$= -\rho \overline{\left(-l'\frac{\partial u}{\partial z}\right)} \left(-l'\frac{\partial u}{\partial z}\right) = -\rho \overline{l'^{2}} \left|\frac{\partial u}{\partial z}\right|^{2}$$

$$= K_{m} \frac{\partial u}{\partial z}, \text{ where } K_{m} = \rho \overline{l'^{2}} \left|\frac{\partial u}{\partial z}\right|^{2}$$
Define $u_{*} = \sqrt{\frac{\tau_{x}}{\rho}} = \left(\overline{u'w'}\right)^{1/2} \text{ then}$

$$\frac{\tau_{x}}{\rho} = \frac{K_{m}}{\rho} \frac{\partial u}{\partial z} = u_{*}^{2}$$

- Momentum flux (surface stress) is proportional to the square of the product of the wind speed gradient (shear) and the turbulent length scale
- Define an "eddy viscosity" or "eddy diffusivity" K_m which is analogous to molecular diffusivity
- Define a velocity scale
 u_{*} for the turbulent
 eddies near the
 surface, called the
 friction velocity

Surface Layer (cont'd)

$$\frac{\tau_x}{\rho} = \frac{K_m}{\rho} \frac{\partial u}{\partial z} = u_*^2$$

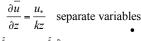
$$K_m = \rho \overline{U^2} \left| \frac{\partial u}{\partial z} \right| = \rho k^2 z^2 \left| \frac{\partial u}{\partial z} \right|$$

$$\frac{K_m}{\rho} \frac{\partial u}{\partial z} = u_*^2 = \frac{\rho k^2 z^2}{\rho} \frac{\left| \frac{\partial u}{\partial z} \right|}{\rho} \frac{\partial u}{\partial z}$$

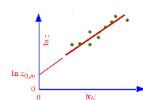
$$\left(\frac{\partial u}{\partial z}\right) = \frac{u_*^2}{k^2 z^2}$$
$$\frac{\partial u}{\partial z} = \frac{u_*}{u_*}$$

- Near the surface, eddies are limited in size by the proximity of the ground, so l' in K_m is l'(z)
- Assume l' = kz, where k ~ 0.4 is an empirical coefficient known as "von Karman's constant"
- Leads to a characteristic relationship for variation of mean wind speed with height: the log-wind profile

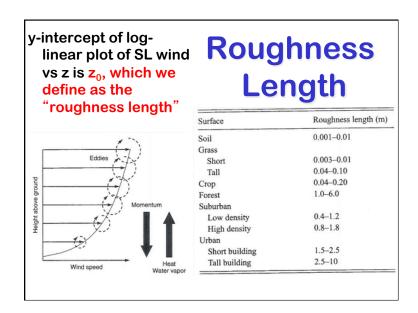
Log Wind Profile

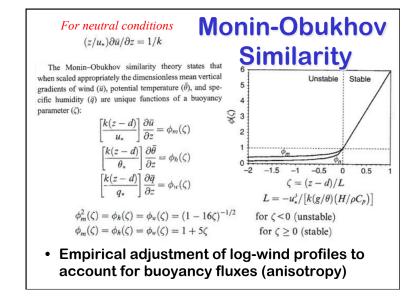


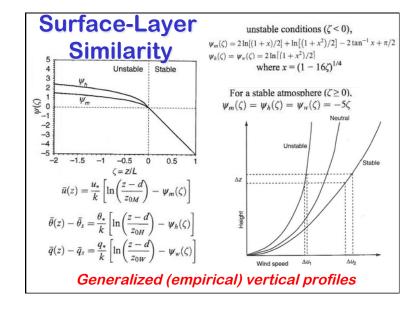
$$\frac{\overline{u}(z) = \frac{u_*}{u} \ln \frac{z}{z}$$



- Mean wind speed in the surface layer is decelerated by friction whose influence is felt aloft through eddy momentum flux
- Varies logarithmically with height
- Y-intercept of log-linear plot of SL wind vs z is z₀, which we define as the "roughness length"







Estimation of Turbulent Fluxes

where
$$\tau = \rho K_M(\partial \bar{u}/\partial z)$$

$$H = -\rho C_p K_H(\partial \bar{\theta}/\partial z)$$

$$E = -\rho K_W(\partial \bar{q}/\partial z)$$

$$K_M = k u_*(z-d)/\phi_m(\zeta)$$

$$K_H = k u_*(z-d)/\phi_h(\zeta)$$

$$K_W = k u_*(z-d)/\phi_w(\zeta)$$

- Fluxes are driven by gradients in u, T, and q
- Fluxes are proportional to friction velocity
- These are simply definitions of K_M , K_H , K_W
- Ohm's Law combined with Similarity

